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Transition State Theory. I.

Transition State Theory = Activated Complex Theory = Absolute Rate Theory

$$\mathrm{H_2} + \mathrm{F} \ \rightleftarrows \ [\mathrm{H_2F}]^\ddagger \overset{k}{\longrightarrow} \mathrm{HF} + \mathrm{H}$$

Assume equilibrium between reactants $H_2 + F$ and the transition state.

$$K^{\ddagger} = \frac{[H_2F]^{\ddagger}}{[H_2][F]}$$

Treat the transition state as a molecule with structure that decays unimolecularly with rate constant k.

$$\frac{d[HF]}{dt} = k [H_2F]^{\ddagger} = kK^{\ddagger} [H_2][F]$$

k has units of sec⁻¹ (unimolecular decay). The motion along the reaction coordinate looks like an antisymmetric vibration of H_2F^{\ddagger} , one-half cycle of this vibration. Therefore k can be approximated by the frequency of the antisymmetric vibration v[sec⁻¹]

 $k \approx v \equiv$ frequency of antisymmetric vibration (bond formation and cleavage looks like antisymmetric vibration)

$$\frac{d[HF]}{dt} = \nu K^{\ddagger} [H2] [F]$$

$$\frac{d[HF]}{dt} = \nu \left[\frac{(q^{\ddagger*} / N)}{(q^{\ddagger+2} / N)(q^{F} / N)} \right] e^{-E^{\ddagger} / kT} [H_{2}] [F]$$

$$\mathbf{K}^{\ddagger} = \left[\frac{\left(\mathbf{q}_{trans}^{\ddagger} / \mathbf{N}\right)}{\left(\mathbf{q}_{trans}^{H_2} / \mathbf{N}\right)}\right] \left(\frac{\mathbf{q}_{rot}^{\ddagger}}{\mathbf{q}_{rot}^{H_2}}\right) \left(\frac{\mathbf{q}_{vib}^{\ddagger}}{\mathbf{q}_{vib}^{\ast H_2}}\right) \left(\frac{\mathbf{g}_0^{\ddagger}}{\mathbf{g}_0^{H_2} \mathbf{g}_0^{F}}\right) \mathbf{e}^{-\mathbf{E}^{\ddagger} / \mathbf{k}T}$$

Reaction coordinate is antisymmetric vibrational mode of H_2F^{\ddagger} . This vibration is fully excited (high T limit) because it leads to the cleavage of the H–H bond and the formation of the H–F bond. For a fully excited vibration

The vibrational partition function for the antisymmetric mode is

$$q_{\nu ib}^{*asym} = \frac{1}{1 - e^{-h\nu/kT}} \quad ' \quad \frac{kT}{h\nu} \qquad since \; e^{-h\nu/kT} \approx 1 - h\nu/kT$$

Note that this is an incredibly important simplification. The unknown v simply disappears! We do not need to estimate it!

$$K^{\ddagger} = \frac{kT}{h\nu} \left[\frac{q_{\text{trans}}^{\ddagger}/N}{\left(q_{\text{trans}}^{\text{H}_{2}}/N\right)\left(q_{\text{trans}}^{\text{F}}/N\right)} \right] \left(\frac{q_{\text{rot}}^{\ddagger}}{q_{\text{rot}}^{\text{H}_{2}}}\right) \left(\frac{q_{\text{vib}}^{\ddagger^{*}}}{q_{\text{vib}}^{\text{H}_{2}}}\right) \left(\frac{g_{0}^{\ddagger}}{g_{0}^{\text{H}_{2}}g_{0}^{\text{F}}}\right) e^{-E^{\ddagger}/kT}$$

where

$$q_{vib}^{\ddagger *'} = \prod_{i=1}^{3n-5-1} \frac{1}{1 - e^{-hv_i/kT}}$$

if transition state is linear

or

$$q_{vib}^{\ddagger*'} = \prod_{i=1}^{3n-6-1} \frac{1}{1 - e^{-hv_i/kT}}$$
 if transiti

f transition state is nonlinear

n is # of atoms in transition state

 $q_{vib}^{\ddagger} \equiv$ partition function from which the antisymmetric vibrational mode is excluded; it has become the *reaction coordinate*

So
$$K^{\ddagger} = K^{\ddagger} = \frac{kT}{hv}K^{\ddagger'}$$
 $K^{\ddagger'} =$ "special" modification of K^{\ddagger}
that excludes the partition function for the

that excludes the partition function for the antisymmetric vibrational mode

What is E[‡]?



Since a molecule cannot have a vibrational energy lower than its zero point energy, the effective barrier along the reaction coordinate is

Lecture #33

$$E^{\ddagger} = V_0 + (ZPE)_{TS} - (ZPE)_R$$

 V_0 is the potential energy difference between the bare barrier (saddle point) and the reactant bare minimum.

For linear H_2F^{\ddagger} , n = 3, so 3n-5-1 = 3 regular vibrational modes, thus

$$E^{\ddagger} = V_0 + \underbrace{\frac{1}{2}h \left[v_{sym.st.}^{\ddagger} + 2v_{bend}^{\ddagger} - v_{H_2} \right]}_{\text{Difference in ZPE}}$$

FORMULATION of k^{TST}

$$\frac{\mathrm{d}[\mathrm{HF}]}{\mathrm{d}t} = v \frac{\mathrm{kT}}{\mathrm{h}v} \mathrm{K}^{\sharp'}[\mathrm{H}_{2}][\mathrm{F}] = \frac{\mathrm{kT}}{\mathrm{h}} \mathrm{K}^{\sharp'}[\mathrm{H}_{2}][\mathrm{F}] = \mathrm{k}^{\mathrm{TST}}[\mathrm{H}_{2}][\mathrm{F}]$$

so

$$\mathbf{k}^{\mathrm{TST}} = \frac{\mathbf{k}\mathrm{T}}{\mathbf{h}}\,\mathbf{K}^{\ddagger'}$$

but not all reactant molecules make it all the way to products — some are reflected back to separated reactants.

Thus,

$$k^{TST} = \kappa \frac{kT}{h} K^{\ddagger'}$$
 where $\kappa \equiv$ transmission coefficient (a fudge factor)

EVALUATION OF k^{TST}

POTENTIAL ENERGY SURFACE KNOWN:

- E^{\ddagger} directly from potential energy surface
- I^{\ddagger} (moment of inertia of transition state) calculate from geometric structure of transition state
- v^{\ddagger} analyze shape of potential in saddle point region
- κ trajectory calculations consider κ = 1 for now.

Lecture #33

$$H_2 + F \rightarrow HF + H$$
 $T = 300K$
 $m_{H_2} = 2$ $m_F = 19$

Translational part

$$\left[\frac{\left(q_{\text{trans}}^{\ddagger}/N\right)}{\left(q_{\text{trans}}^{\text{H}_{2}}/N\right)\left(q_{\text{trans}}^{\text{F}}/N\right)}\right] = \frac{Nh^{3}}{(2\pi kT)^{3/2}} \left(\frac{m^{\ddagger}}{m_{\text{H}_{2}}m_{\text{F}}}\right)^{3/2} = 6 \times 10^{23} \,\text{mol}^{-1} \frac{(6.63 \times 10^{-34} \,\text{J} \cdot \text{s})^{3}}{(2\pi \times 1.38 \times 10^{-23} \,\text{J/K} \times 300 \,\text{K})^{3/2}} \left(\frac{6 \times 10^{23} \times 0.021}{0.002 \times 0.019 \,\text{kg}}\right)^{3/2} = 2.52 \times 10^{-7} \,\text{m}^{3} \,\text{mol}^{-1}$$

Rotational part

$$\sigma_{H_2} = 2 \qquad \sigma^{\ddagger} = 1 \qquad q_{rot} = \frac{8\pi^2 IkT}{\sigma h^2}$$

$$I_{H_2} = 4.56 \times 10^{-48} m^2 kg \qquad I^{\ddagger} = 1.24 \times 10^{-46} m^2 kg \qquad (assume linear transition state)$$

$$\frac{q_{rot}^{\ddagger}}{q_{rot}^{H_2}} = \frac{I^{\ddagger}}{I_{H_2}} \frac{\sigma_{H_2}}{\sigma^{\ddagger}} = 54.4$$

Vibrational Part

 H_2F^{\ddagger} is a linear transition state (assumed)

3n-5-1 = 3 vibrational degrees of freedom (one vibration is reaction coordinate)

$$\frac{hv_s^{\ddagger}}{k} = 5771 \text{K}$$
 stretch $\frac{hv_b^{\ddagger}}{k} = 573 \text{K}$ (doubly degenerate) bend

 $hv_{H_2}/k = 6323K$

$$\frac{q_{vib}^{\ddagger s'}}{q_{vib}^{\ast H_2}} = \frac{\left(1 - e^{-hv_s^{\ddagger}/kT}\right)^{-1} \left(1 - e^{-hv_b^{\ddagger}/kT}\right)^{-2}}{\left(1 - e^{-hvH_2/kT}\right)^{-1}} = 1.38$$

Lecture #33

Electronic part

$$g_{0}^{\ddagger} = 2(S = 1/2) \quad g_{0}^{F} = 6(L = 1, S = 1/2) \quad g_{0}^{H_{2}} = 1$$

(spin orbit splitting of F ²P_{1/2} - ²P_{3/2} is 404 cm⁻¹)
$$\frac{g_{0}^{\ddagger}}{g_{0}^{F}g_{0}^{H_{2}}} = \frac{1}{3}$$

Calculate E[‡]

$$V_{0} = 3.8 \text{ kJmol}^{-1} \qquad v_{s}^{\ddagger} = 1.20 \times 10^{14} \text{ s}^{-1} \qquad v_{b}^{\ddagger} = 1.19 \times 10^{13} \text{ s}^{-1} \text{ (reasonable guesses)}$$
(How do we guess values for v_{s}^{\ddagger} and v_{b}^{\ddagger} ?)
$$v_{H_{2}} = 1.32 \times 10^{14} \text{ s}^{-1}$$

$$E^{\ddagger} = V_0 + \frac{1}{2} hN \left[v_s^{\ddagger} + 2v_b^{\ddagger} - v_{H_2} \right] = 6.1 \text{ kJ mol}^{-1}$$

Calculate kT/h

$$kT/h = \frac{1.38 \times 10^{-23} \text{ J/K} \times 300 \text{ K}}{6.63 \times 10^{-34} \text{ J} \cdot \text{s}} = 6.24 \times 10^{12} \text{ s}^{-1}$$

Putting it all together:

$$k^{\text{TST}} = \kappa \frac{kT}{h} K^{\pm'}$$

= 1(6.24 × 10^{12} s^{-1})(2.52 × 10^{-7} m^{3} \text{mol}^{-1})(54.4)(1.38) \frac{1}{3} e^{-6.1/\text{RT}}
= 3.93 × 10⁷ e^{-6.1/RT}
$$k^{\text{TST}} = 3.40 × 10^{6} m^{3} \text{mol}^{-1} \text{s}^{-1} \text{ at } 300 \text{K}$$

 $k^{\text{EXP}} = 2.70 \text{ x } 10^6 \text{ m}^3 \text{ mol}^{-1} \text{ s}^{-1}$ acceptable agreement

Experimental value is smaller because κ is probably not 1. Sometimes k^{TST} will be smaller than k^{EXP} because of tunneling. This model for k^{TST} does not take the quantum mechanical phenomenon of tunneling into account. Tunneling can make the reaction rate become faster than the k^{TST} prediction.

If $k^{TST} < k^{EXP}$, it may mean that there is some tunneling contribution.